

Geometric data analysis

Lecture 5/7 – Riemannian metrics and geodesics

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Thursday, 9am–12pm – 7 lectures

Faculté de médecine, Hôpital Cochin, rooms 2001 + 2005

Validation: project + quizz

Recap of the previous lectures

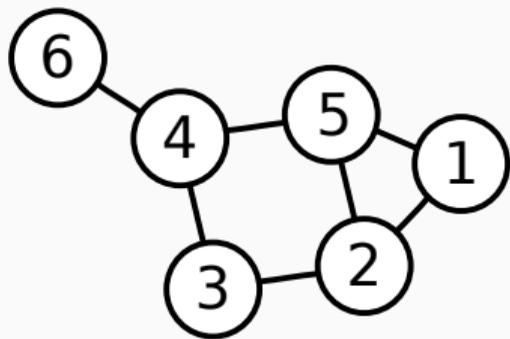
To mitigate the **curse of dimensionality**, we use:

- **Expert** knowledge: high-quality features.
- Relevant **families** of functions: kernels, convolutional networks.
- Relevant **neighborhood** structures: graphs.

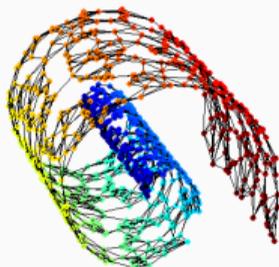
Main challenge: **local** implementation \implies **global** understanding.

Produce **guidelines** and **insights** for practitioners.

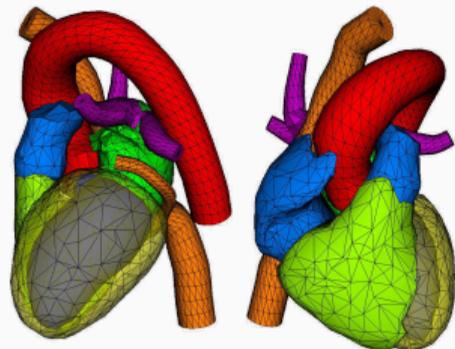
Large graphs are best understood as continuous objects [Pey11, EPW11]



Simple graph.

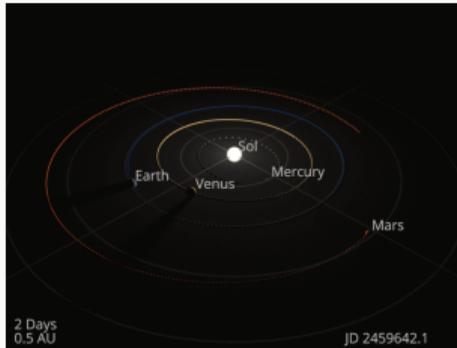


Manifold **hypothesis**.

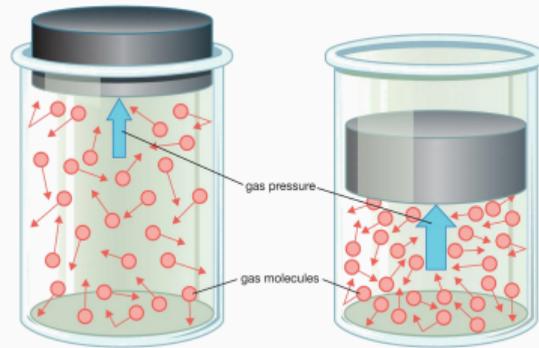


Physical manifold.

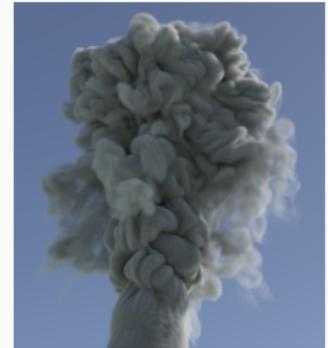
Long history in physics [Dat18, Bria, NWRC22]



The **Solar** system.



The **ideal gas** model.



Fluid simulation.

Lecture 5 – From discrete graphs to **continuous spaces**:

- The Poincaré disk.
- Local metrics and geodesics.

Lecture 6 – From discrete samples to **continuous distributions**:

- Duality: distributions and adversarial norms.
- Information geometry, kernels and optimal transport.

Lecture 7 – **Hardware** bottlenecks:

- Registers, parallel cores and compilers.
- Current trends.

Textbooks and introductions – in English:

- *Poincaré and his disk* – Étienne Ghys, 2006.
- *Hyperbolic geometry* – Cannon et al., 1997.
- *Riemannian geometry: an introduction to curvature* – John M. Lee, 1997.
- *Geometric data analysis, beyond convolutions* – my PhD thesis, 2020.

Lecture notes available on my website – in French:

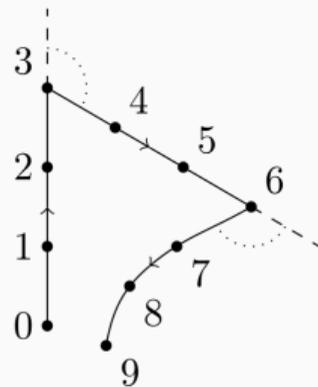
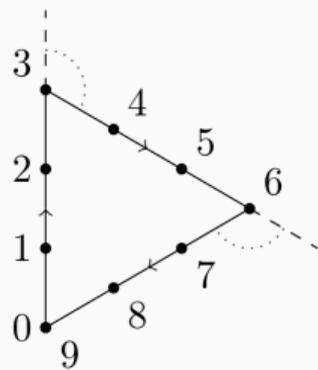
- *Culture mathématique.*
- *Introduction à la géométrie riemannienne par l'étude des espaces de formes.*

The Poincaré disk

Science and hypothesis – Henri Poincaré, 1902

The Non-Euclidean World. – If geometrical space were a framework imposed on **each** of our representations considered **individually**, it would be **impossible** to represent to ourselves an image without this framework, and we should be quite **unable** to change our geometry.

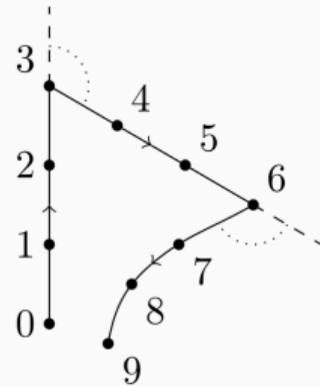
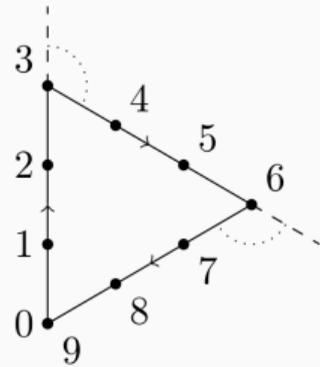
But this is not the case; geometry is only the summary of the laws by which these images succeed each other.



Science and hypothesis – Henri Poincaré, 1902

There is nothing, therefore, to prevent us from imagining a **series of representations**, similar in every way to our ordinary representations, but succeeding one another according to laws which **differ** from those to which we are accustomed.

We may thus conceive that beings whose education has taken place in a medium in which those laws would be so different, might have **a very different geometry** from ours.

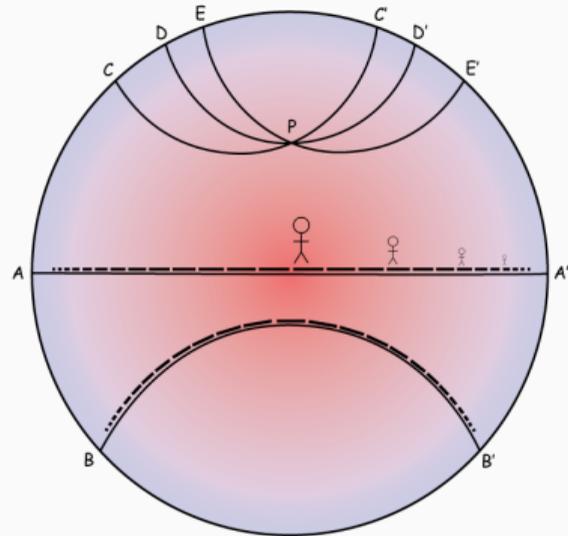


Suppose, for example, a world **enclosed in a large sphere** and subject to the following laws:

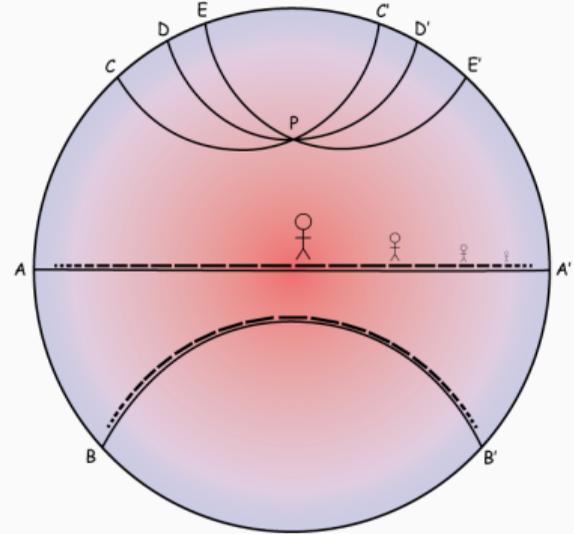
- The **temperature is not uniform**; it is **greatest at the centre**, and gradually decreases as we move towards the **circumference** of the sphere, where it is **absolute zero**.

The **law of this temperature is as follows**:

if R be the radius of the sphere, and r the distance of the point considered from the centre, the absolute temperature will be **proportional** to $R^2 - r^2$.



- Further, I shall suppose that in this world **all bodies have the same coefficient of dilatation**, so that the linear **dilatation** of any body is proportional to its absolute **temperature**.
- Finally, I shall assume that a body transported from one point to another of different temperature is **instantaneously** in thermal **equilibrium** with its new environment.

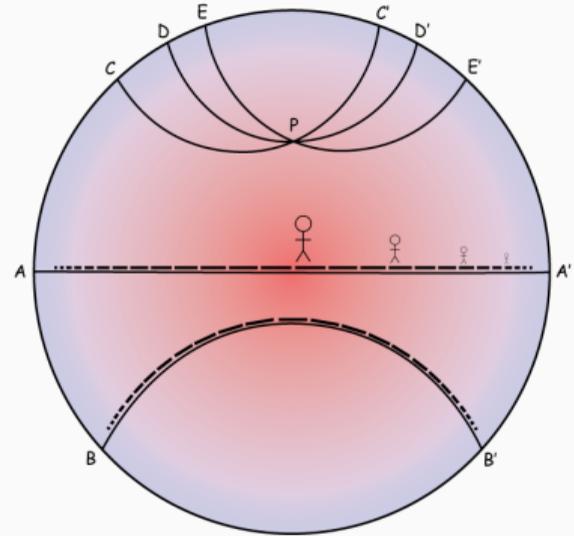


Science and hypothesis – Henri Poincaré, 1902 [Nor]

There is **nothing** in these hypotheses either **contradictory** or **unimaginable**. A moving object will become **smaller and smaller** as it approaches the **circumference** of the sphere.

Let us observe, in the first place, that although from the point of view of **our** ordinary geometry this world is **finite**, **to its inhabitants it will appear infinite**.

As they approach the surface of the sphere they become **colder**, and at the same time smaller and smaller. The steps they take are therefore also **smaller** and smaller, so that they can never reach the **boundary** of the sphere.



The Poincaré metric is locally a Euclidean metric

With $x^2 + y^2 < 1$, we define:

$$d\left((x, y) \rightarrow (x, y) + (dx, dy)\right) = 2 \frac{\sqrt{dx^2 + dy^2}}{1 - (x^2 + y^2)}$$

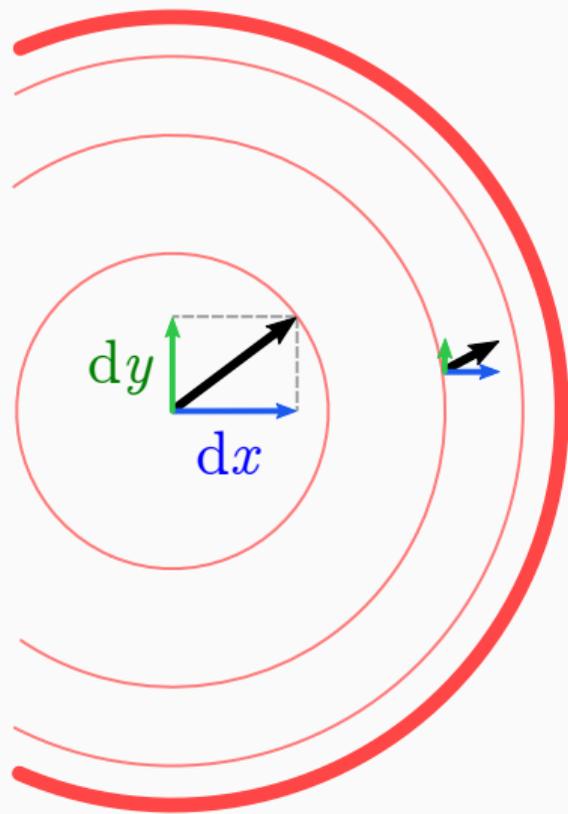
In other words:

$$d^2\left((x, y) \rightarrow (x, y) + (dx, dy)\right) = 4 \frac{dx^2 + dy^2}{(1 - (x^2 + y^2))^2}$$

This **local** Euclidean metric is the **Riemannian** metric:

$$\|(dx, dy)\|_{(x,y)}^2 = \langle (dx, dy), g_{(x,y)}(dx, dy) \rangle$$

$$g_{(x,y)} = \begin{pmatrix} 4/(1 - (x^2 + y^2))^2 & 0 \\ 0 & 4/(1 - (x^2 + y^2))^2 \end{pmatrix}$$



Length of curve

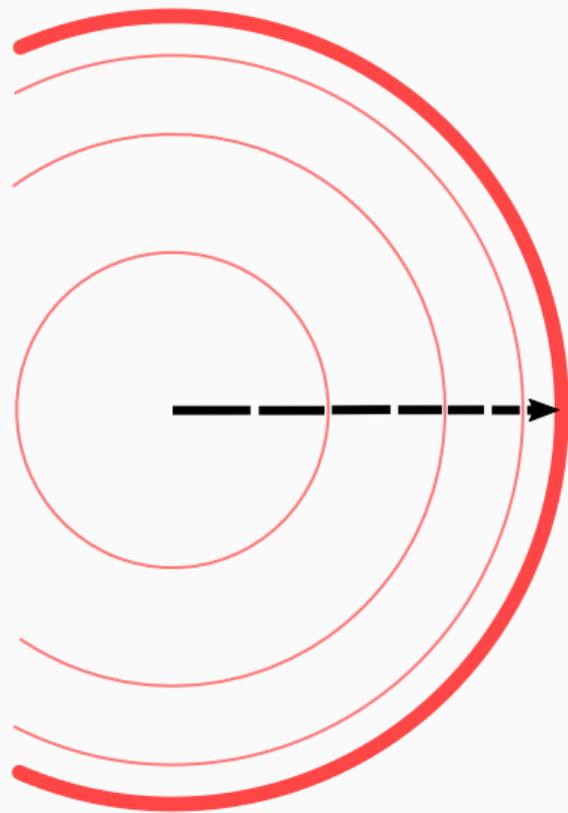
If $\gamma : [0, 1] \rightarrow B(0, 1)$ is a **smooth path**, we define:

$$\ell(\gamma) = \int_0^1 \|\dot{\gamma}(t)\|_{\gamma(t)} dt$$

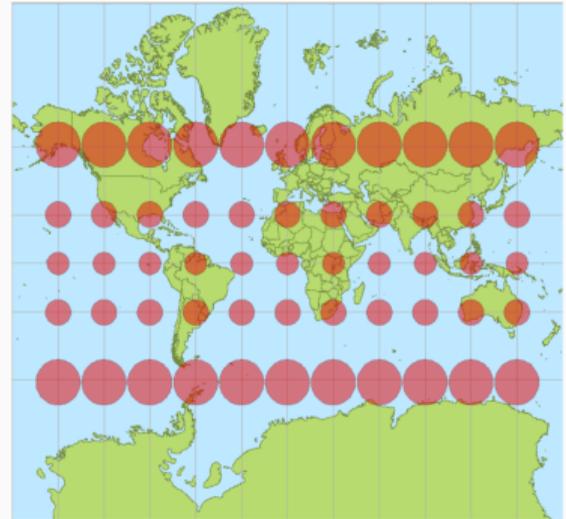
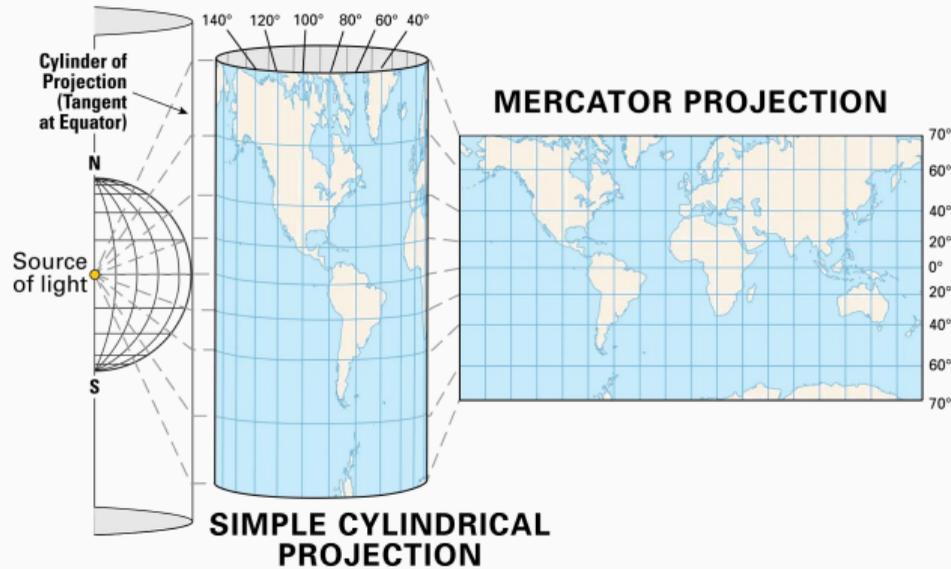
For the **straight path** $\gamma(t) = (t, 0)$, we find:

$$\begin{aligned}\ell(\gamma) &= \int_0^1 \|(1, 0)\|_{(t, 0)} dt \\ &= \int_0^1 \frac{2}{1-t^2} dt \\ &= \int_0^1 \frac{2}{(1+t)(1-t)} dt = +\infty\end{aligned}$$

The Poincaré disk is **a universe in a nutshell**.



Tissot's indicatrix [Brib, Kü04]



Tissot's indicatrix at location (x, y) is a **unit ball** for the local metric.
This ellipsoid allows us to **depict distortions** in cartography
and fully describes a **Riemannian metric** on the 2D map.

Equivalent descriptions of the Poincaré segment $(-1, +1)$

Stereographic projections define bijections between the segment I , the half-circle J and the half-line H .

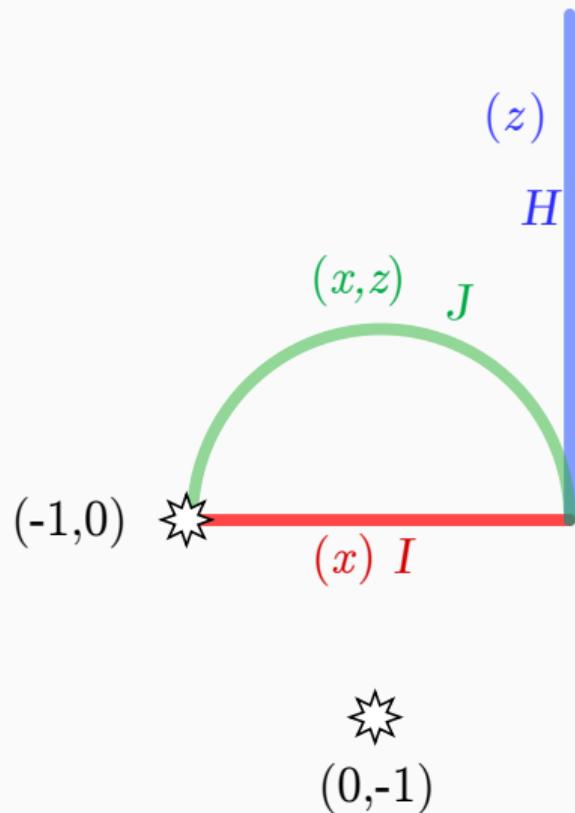
If we endow I with the Poincaré metric:

$$\|(dx)\|_{(x) \in I} = 2 \frac{\sqrt{dx^2}}{1-x^2},$$

then J and H are endowed with:

$$\|(dx, dz)\|_{(x,z) \in J} = \frac{\sqrt{dx^2 + dz^2}}{z}$$

$$\|(dz)\|_{(z) \in H} = \frac{\sqrt{dz^2}}{z}$$



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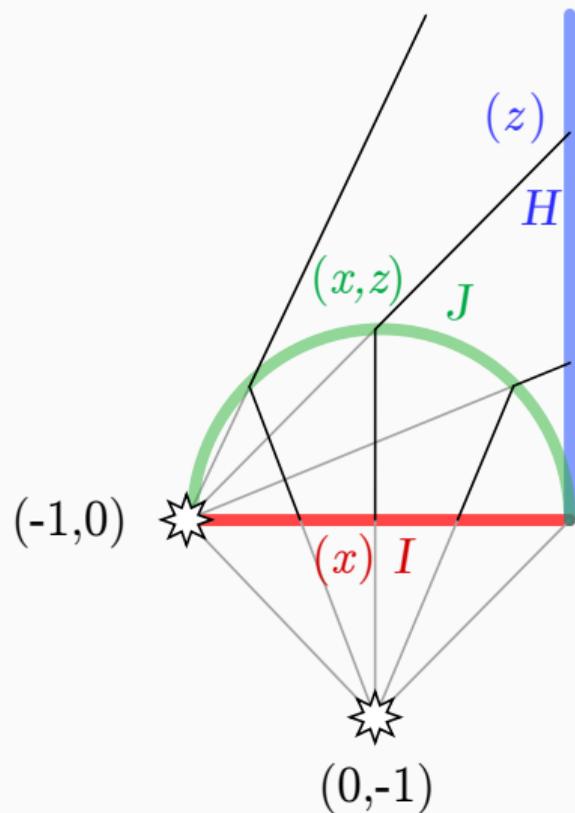
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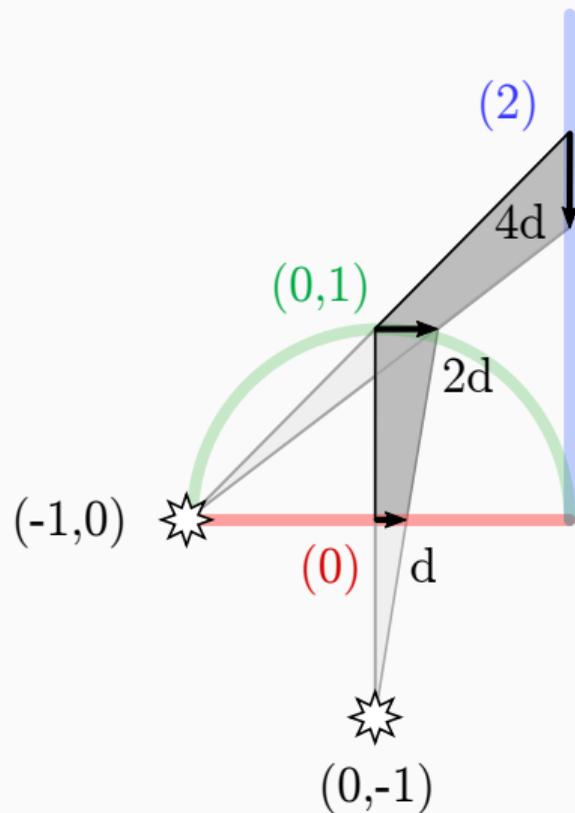
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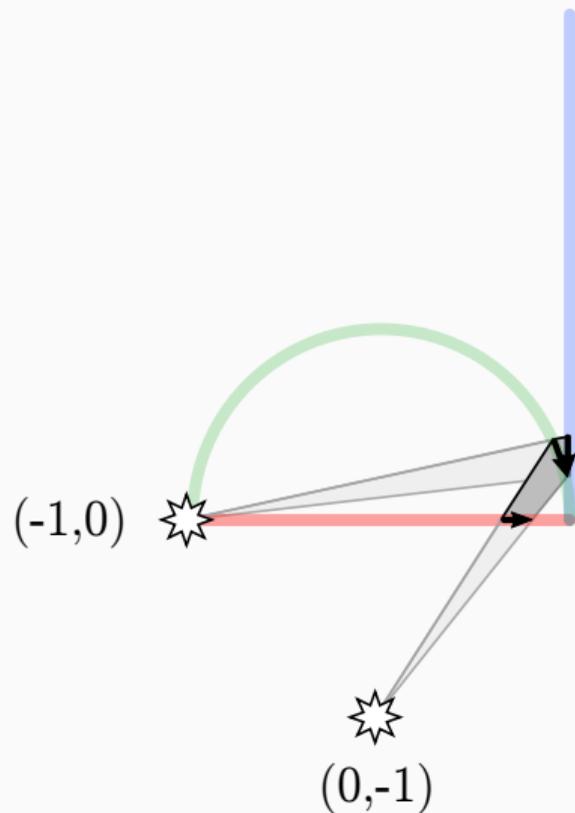
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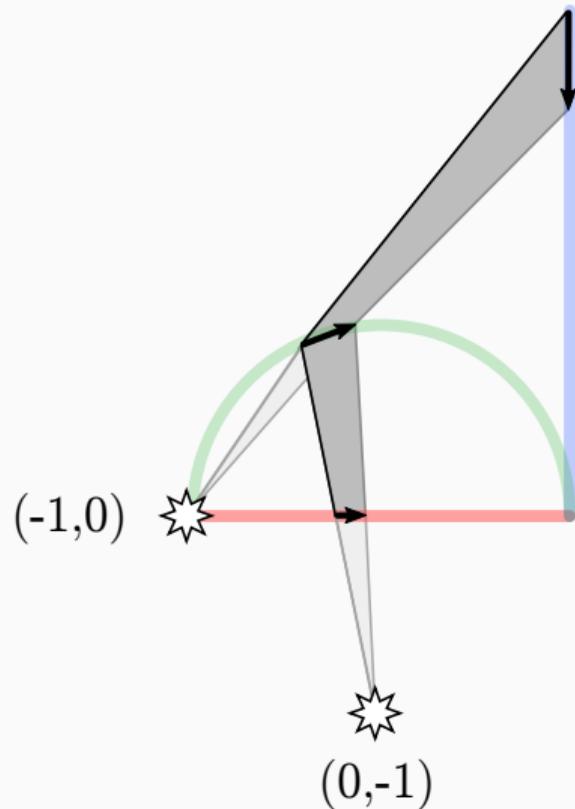
If we endow I with the Poincaré metric:

$$\|(dx)\|_{(x) \in I} = 2 \frac{\sqrt{dx^2}}{1-x^2},$$

then J and H are endowed with:

$$\|(dx, dz)\|_{(x,z) \in J} = \frac{\sqrt{dx^2 + dz^2}}{z}$$

$$\|(dz)\|_{(z) \in H} = \frac{\sqrt{dz^2}}{z}$$



Equivalent descriptions of the Poincaré disk

Stereographic projections define bijections between the disk I , the hemisphere J and the half-plane H .

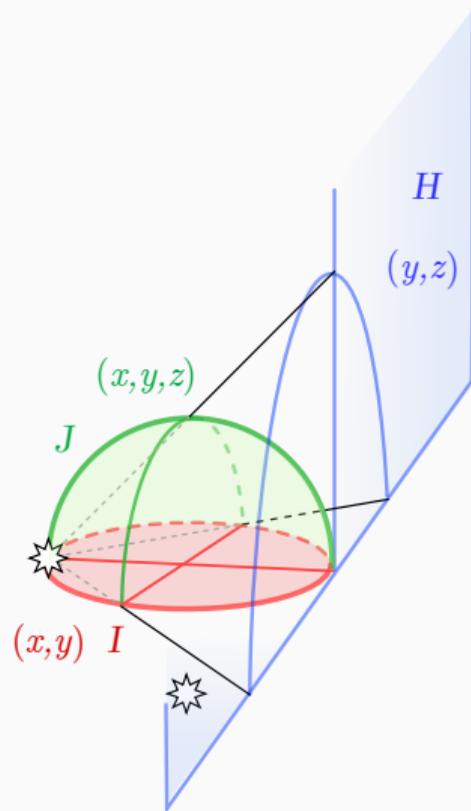
If we endow I with the Poincaré metric:

$$\|(dx, dy)\|_{(x,y) \in I} = 2 \frac{\sqrt{dx^2 + dy^2}}{1 - (x^2 + y^2)},$$

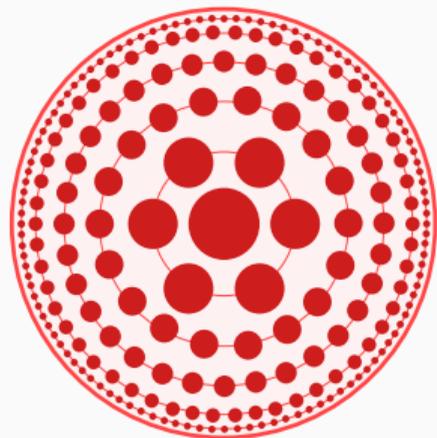
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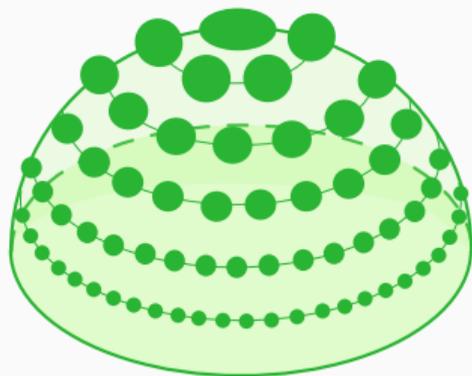
$$\|(dy, dz)\|_{(y,z) \in H} = \frac{\sqrt{dy^2 + dz^2}}{z}$$



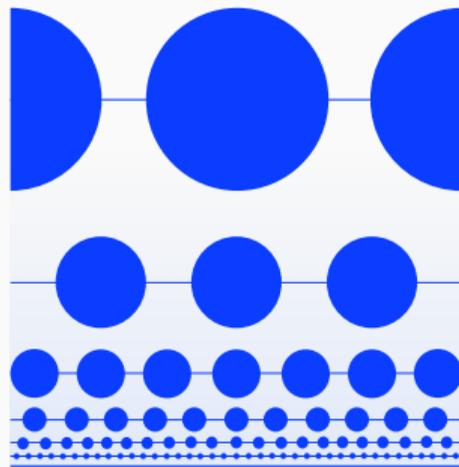
Tissot's indicatrix on the Poincaré disk, hemisphere and half-plane



$$2 \frac{\sqrt{dx^2 + dy^2}}{1 - (x^2 + y^2)}$$

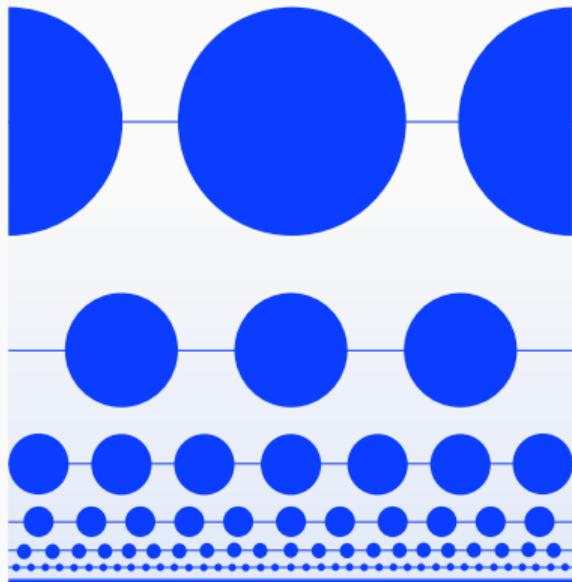


$$\frac{\sqrt{dx^2 + dy^2 + dz^2}}{z}$$

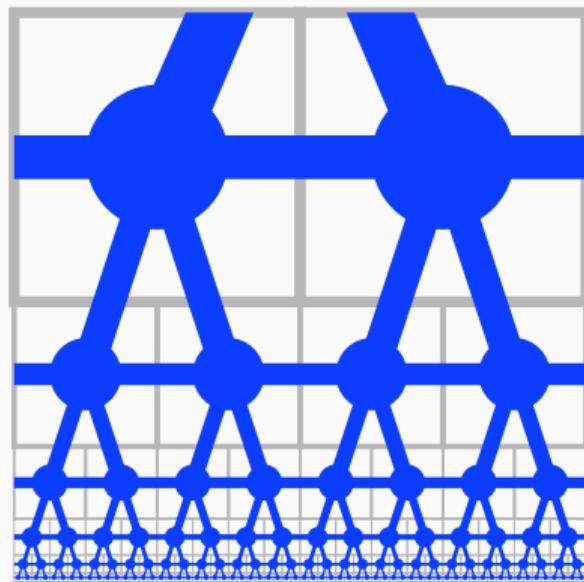


$$\frac{\sqrt{dy^2 + dz^2}}{z}$$

The discrete Poincaré grid



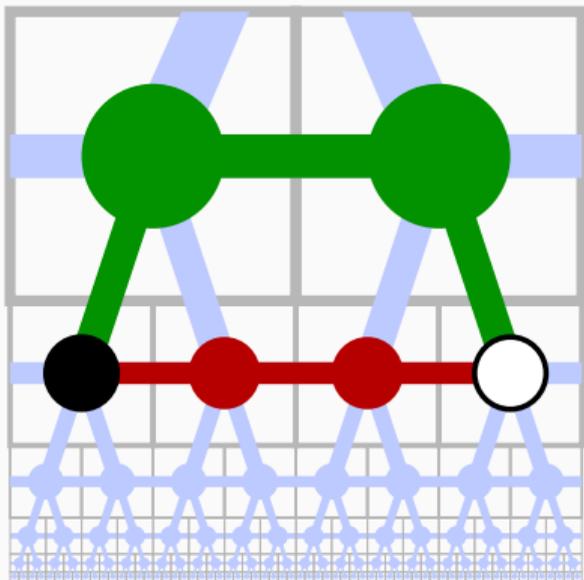
$$\frac{\sqrt{dy^2 + dz^2}}{z}$$



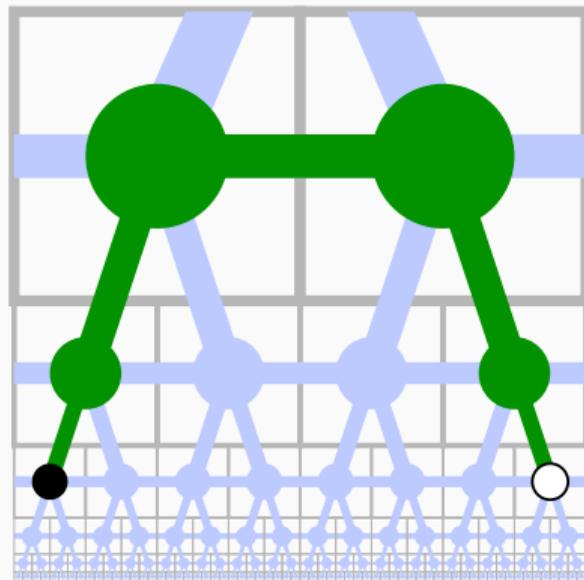
Octave grid based on a dyadic tree.

Geodesics on the Poincaré disk

Geodesics on the Poincaré grid



The **green** and **red** paths
have the same length.

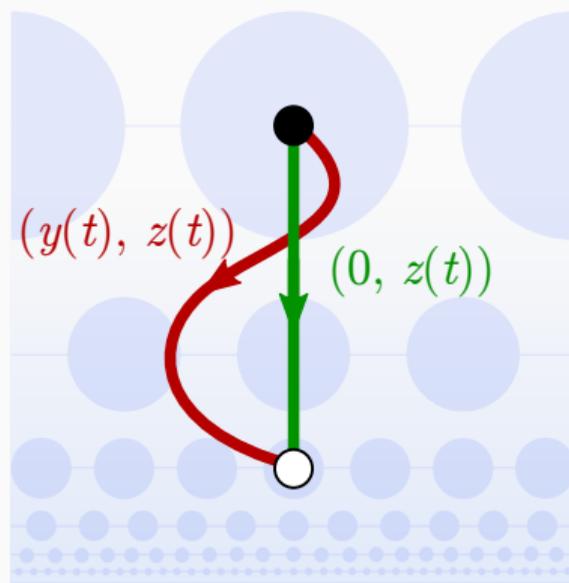


Going **up** is faster than
travelling **sideways**.

Vertical geodesics on the Poincaré half-plane

If the source and target points belong to the **vertical axis** $y = 0$, the shortest path is **straight**:

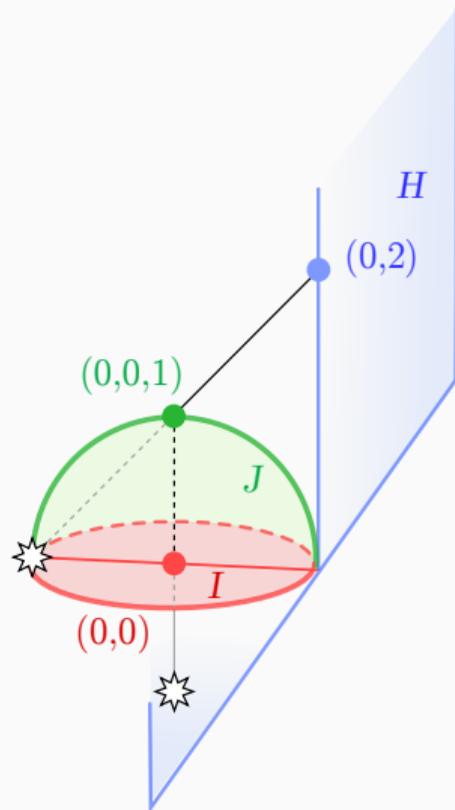
$$\begin{aligned} \ell(y(t), z(t)) &= \int_0^1 \frac{\sqrt{\dot{y}(t)^2 + \dot{z}(t)^2}}{z(t)} dt \\ &\leq \int_0^1 \frac{\sqrt{\dot{z}(t)^2}}{z(t)} dt \\ &= \ell(0, z(t)). \end{aligned}$$



Centers of the three Poincaré models

Using our stereographic projections, the North pole $(x, y, z) = (0, 0, +1)$ of the **hemisphere** corresponds to:

- the center $(x, y) = (0, 0)$ of the **disk**,
- the point $(y, z) = (0, +2)$ of the **half-plane**.



Geodesics on the Poincaré hemisphere

The vertical axis in the **half-plane** is equivalent to the “**Greenwich meridian**” in the **hemisphere**.

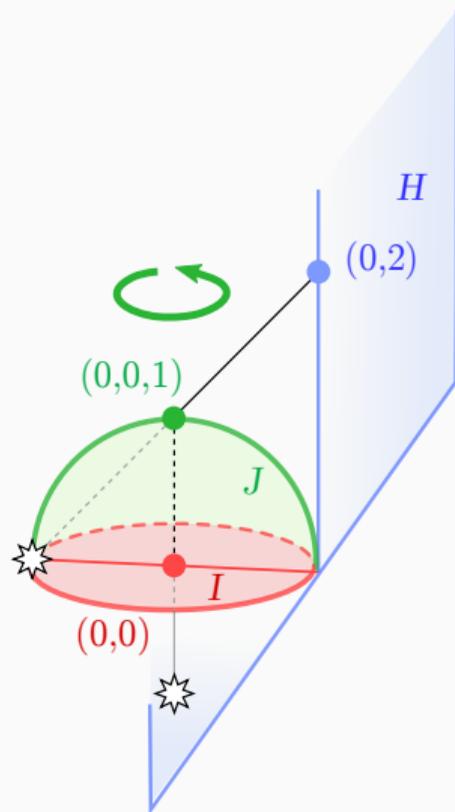
Since the Poincaré metric:

$$\|(dx, dy, dz)\|_{(x,y,z) \in J} = \frac{\sqrt{dx^2 + dy^2 + dz^2}}{z}$$

is invariant by **rotations**:

$$(x, y, z) \mapsto (\cos(\theta) x, \sin(\theta) y, z),$$

all great circles that pass through the North pole are also **geodesic curves**.



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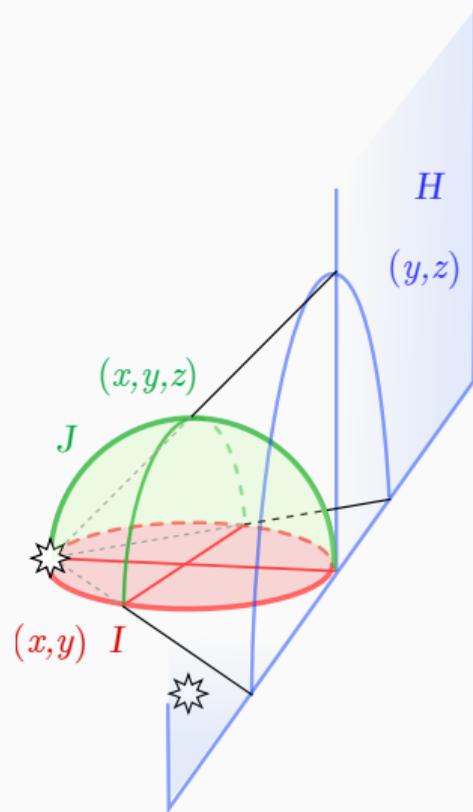
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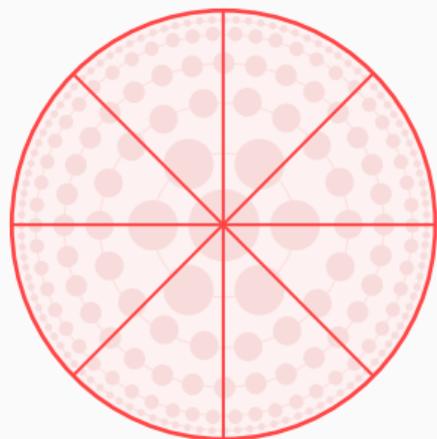
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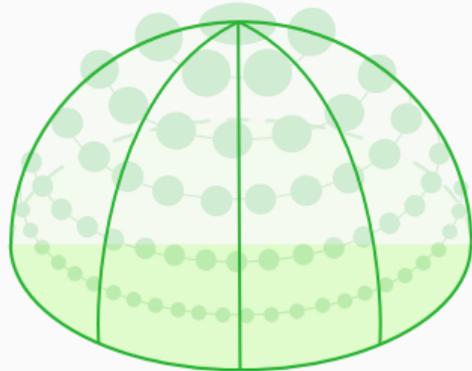
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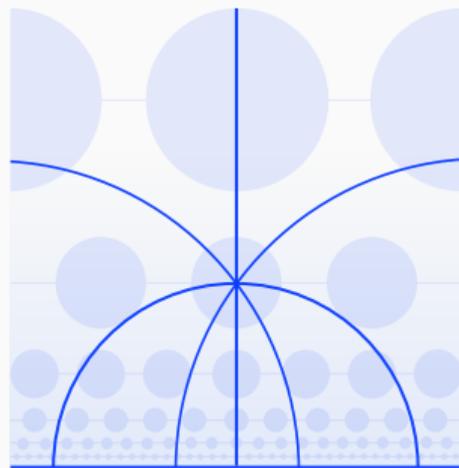
Meridians on the three Poincaré models



Diameters.



Meridians.



Half-circles that pass through $(y, z) = (0, +2)$.

Geodesics on the Poincaré half-plane

Half-circles that are perpendicular to the horizontal axis **are all geodesics**.

This is because the Poincaré metric:

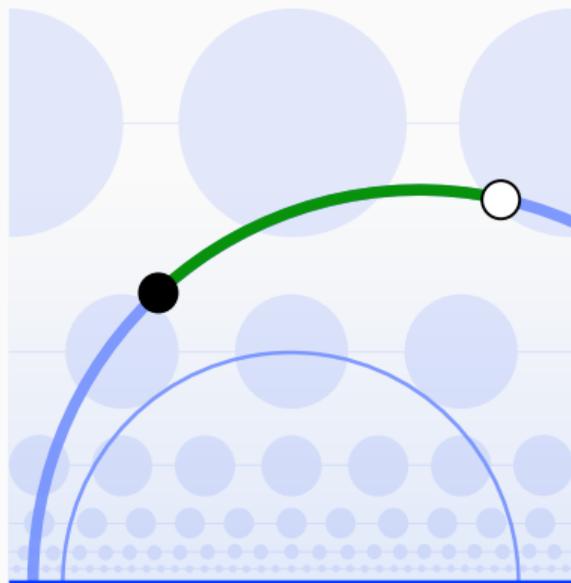
$$\|(dy, dz)\|_{(y,z) \in H} = \frac{\sqrt{dy^2 + dz^2}}{z}$$

is invariant by **horizontal translations**:

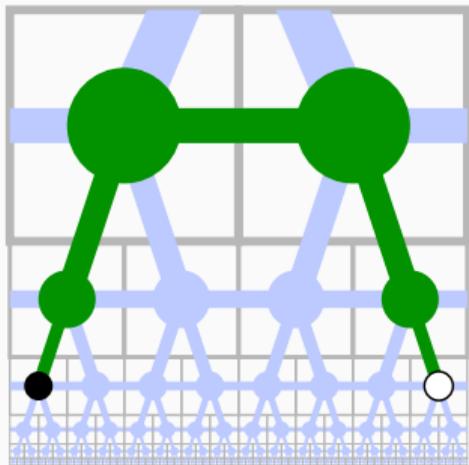
$$(y, z) \mapsto (y + \Delta y, z)$$

and **scalings** with a positive constant $a > 0$:

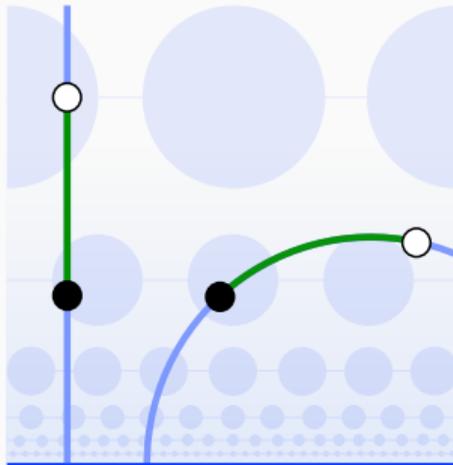
$$(y, z) \mapsto (ay, az) .$$



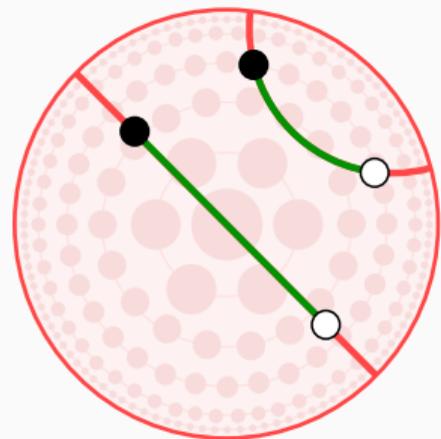
Geodesics on the Poincaré grid, half-plane and disk



Up-and-down paths.



Half-circles and
vertical lines.

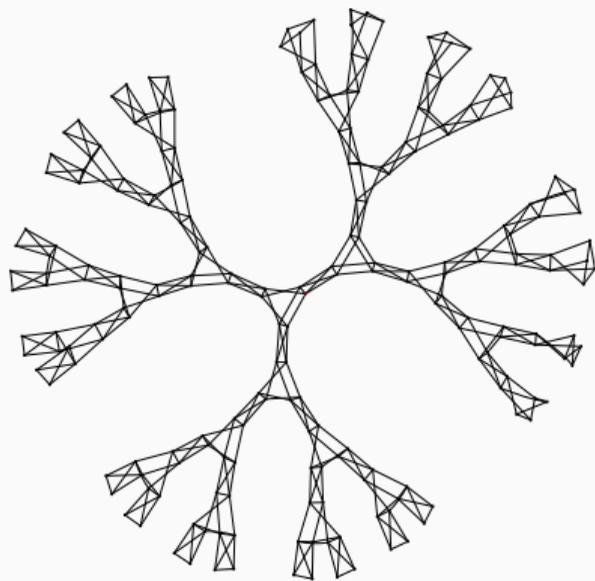


Orthogonal circles and
diameters.

A continuous tree



Circle Limit IV by **M.C. Escher**:
a **regular** tiling of the Poincaré disk.



The Cayley graph of $SL_2(\mathbb{Z})$
is to the **Poincaré** disk what a
regular grid is to the **Euclidean** plane.

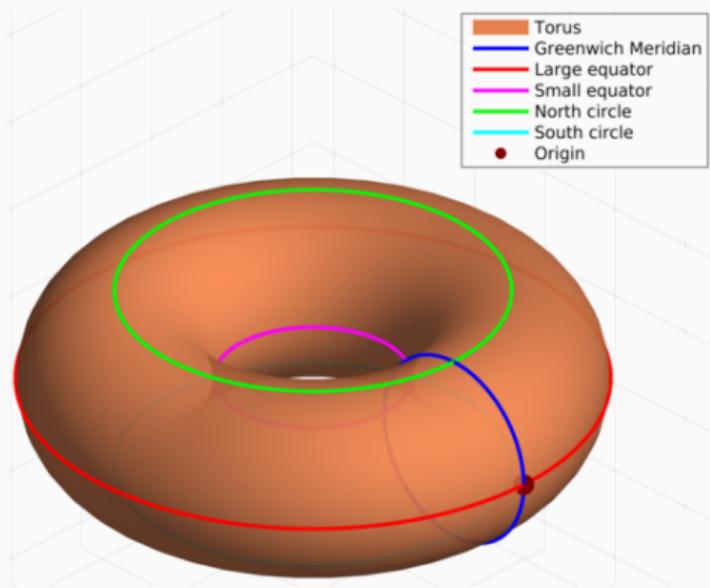
Recap on the Poincaré disk

We should **remember** that:

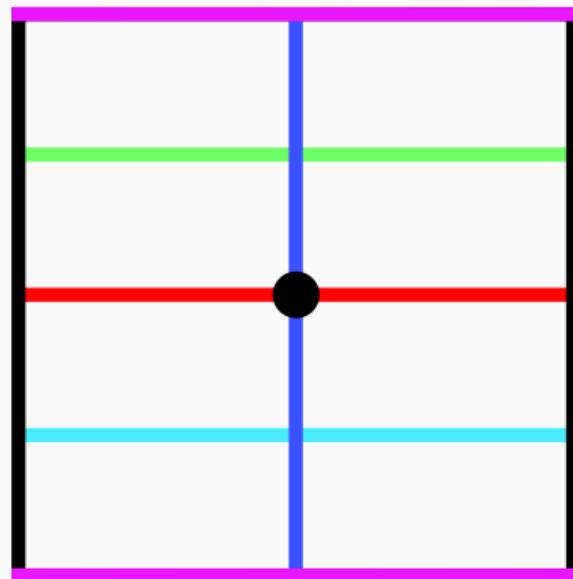
- A **Riemannian** metric is a smooth field of Euclidean norms. It is equivalent to **Tissot's indicatrix** in cartography.
- **Convenient** way of defining arbitrary geometries: the Poincaré disk is a continuous tree.
- **Changes of coordinates** are key to eloquent proofs. Don't get stuck on one parameterization.
- **Discrete** \leftrightarrow **Continuous** analogies go both ways.

Riemannian geometry, in practice

A convenient way of working with surfaces

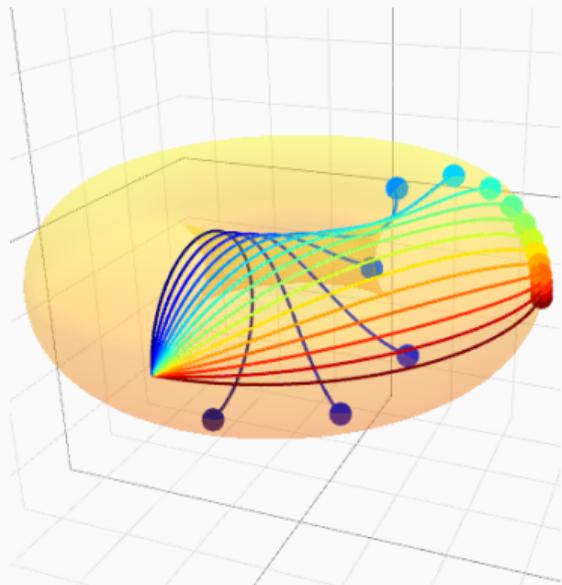


As a surface **embedded in 3D**.

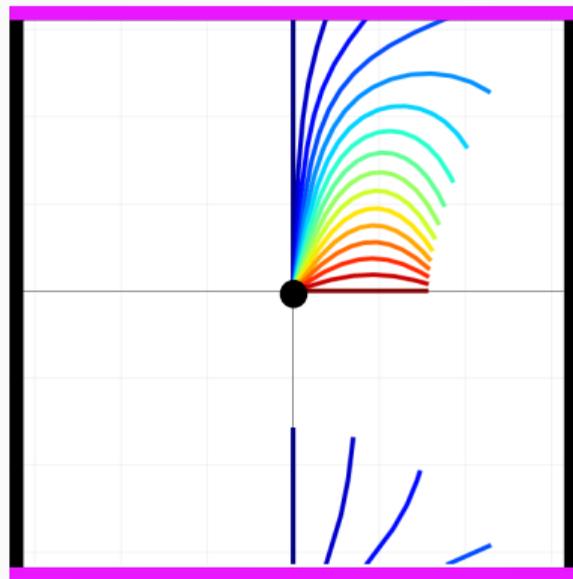


The torus as a **flat 2D square**.

A convenient way of working with surfaces



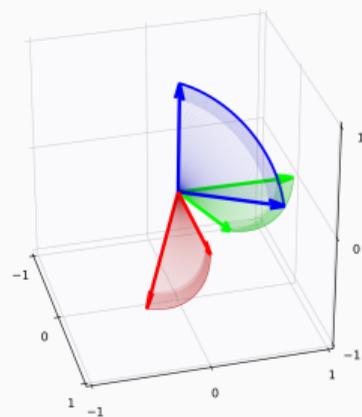
Geodesics on the donut.



The torus as a **curved 2D square**.



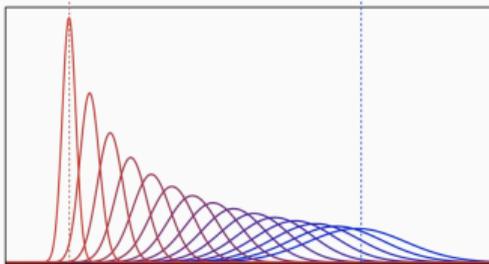
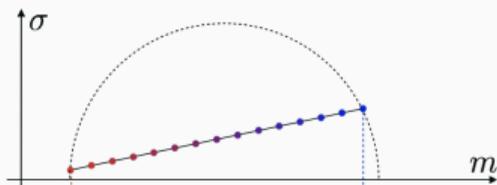
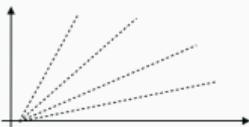
Covariance matrices: geodesics for the
Euclidean, affine-invariant and
log-Euclidean metrics.



3D rotations: geodesics for
the Lie group structure.

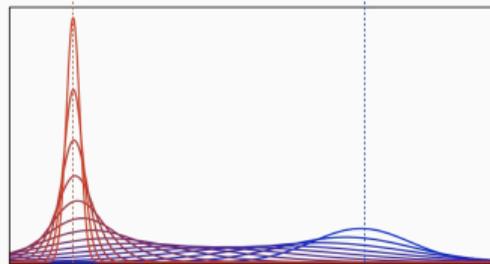
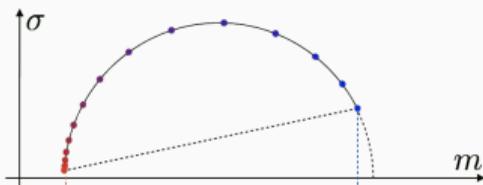
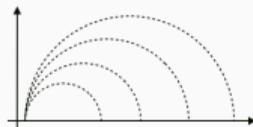
Probability distributions – see Lecture 6! [PC18]

Optimal Transport
(Euclidean)



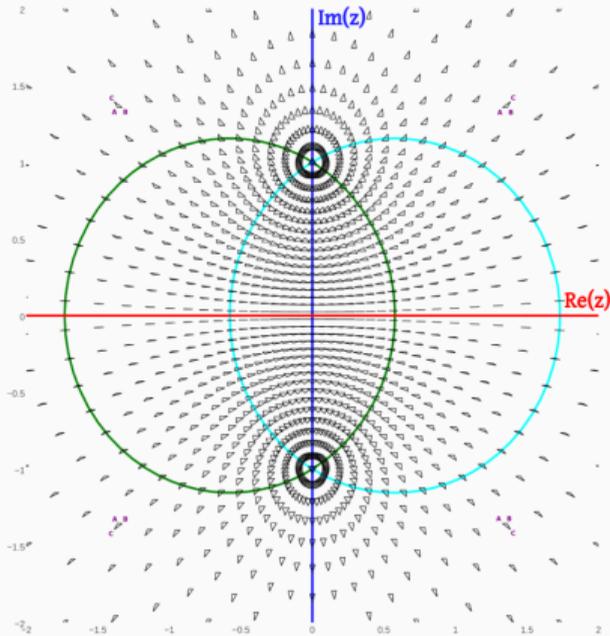
Gaussians + **Wasserstein** metric
= **Euclidean**.

Fisher-Rao
(hyperbolic)

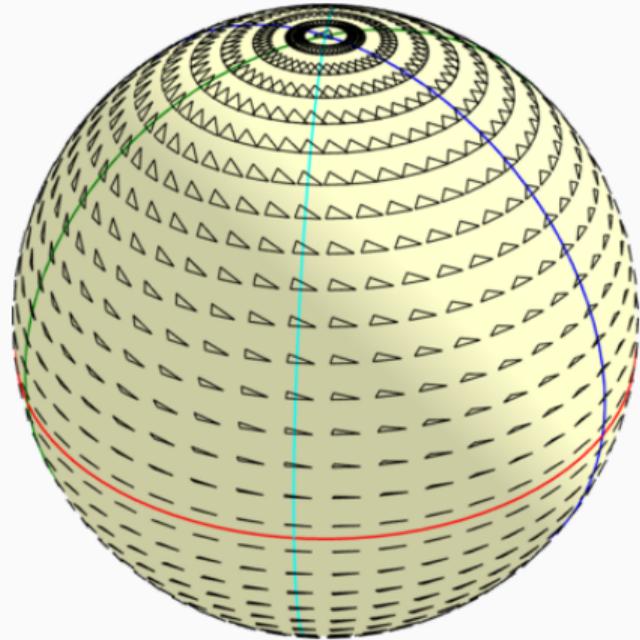


Gaussians + relative **entropy**
= **Poincaré**.

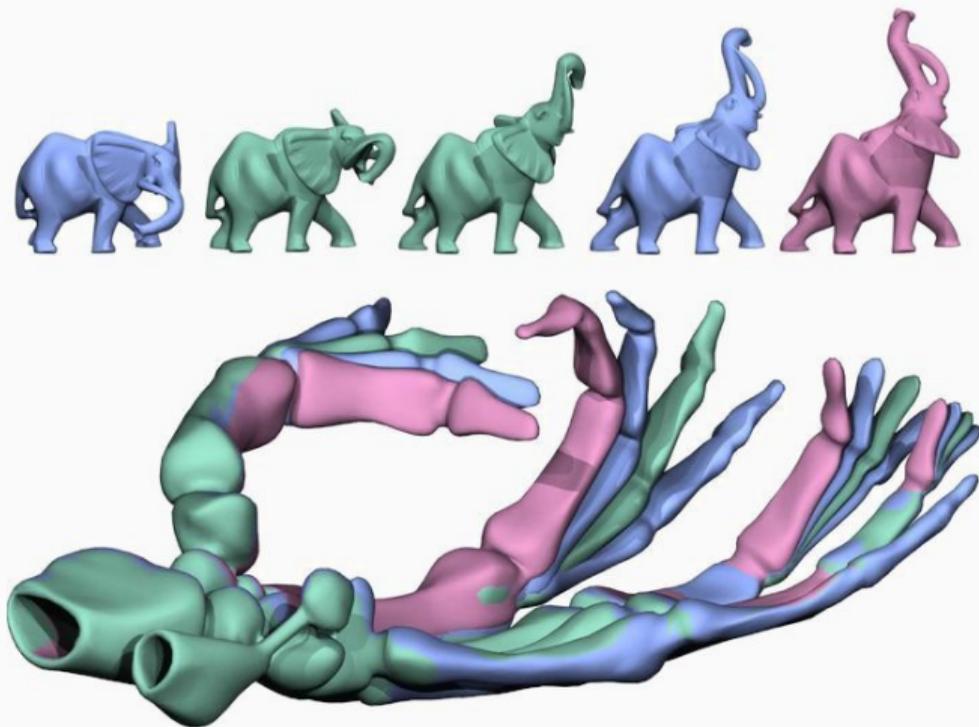
Shape metrics – remember Lecture 1!



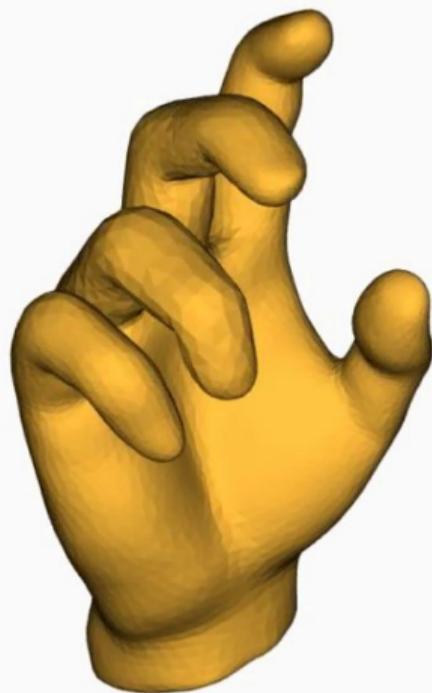
The **plane** of triangle shapes.



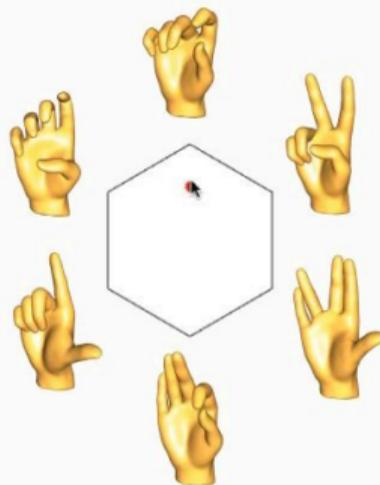
The **sphere** of triangle shapes.



Geodesics in spaces of elephants and skeletons.

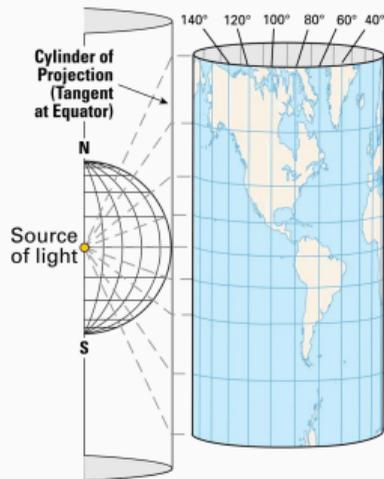


🖥️ screen captured

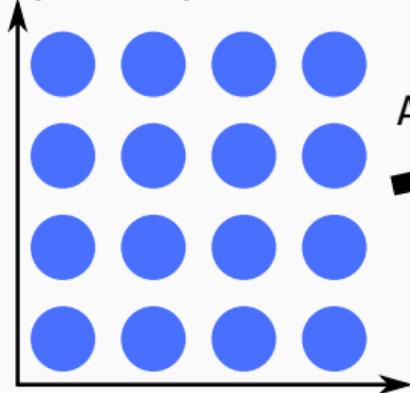


Barycentric interpolation in a space of hands.

Network architectures are “projections” from parameter space to function space



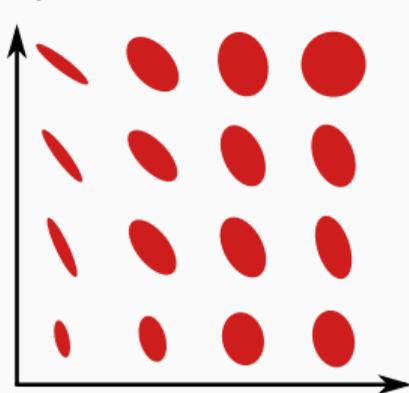
Space of parameters



Architecture



Space of estimators



Standard gradient descent on the parameters

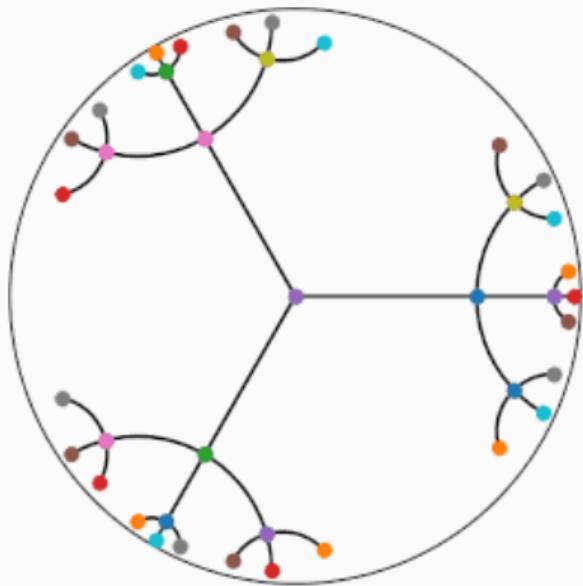


Riemannian gradient descent on the network.

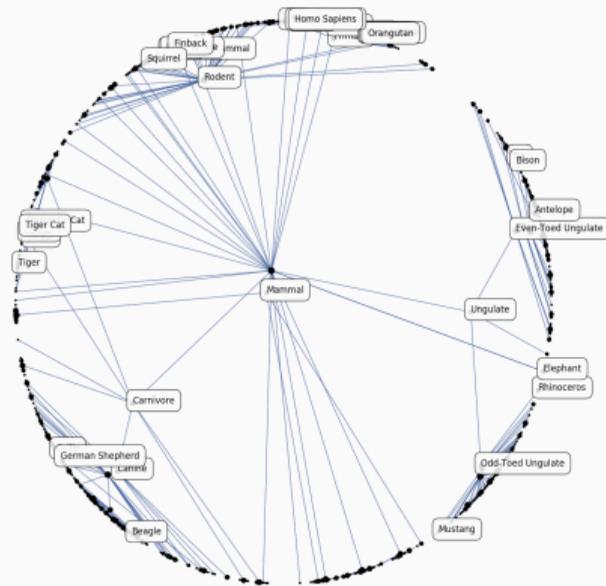
Core idea behind the **natural gradient**, neural **tangent kernels**,

(Wasserstein) **gradient flows**...

Graph and point cloud embeddings [SDSGR18, NK17]



Embedding a **tree** in the Poincaré disk.



WordNet **mammals** subtree

UMAP (Uniform **Manifold** Approximation and Projections)
also works with Riemannian metrics.

The geodesic equation – Hamiltonian form

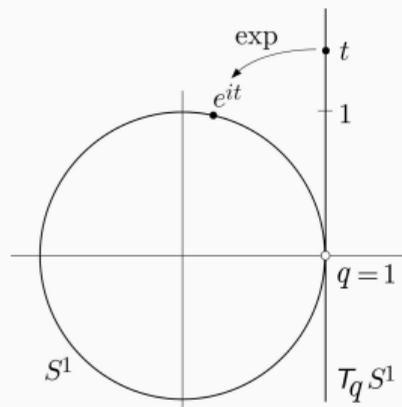
If $H(q, p) = \frac{1}{2} \langle p, g_q^{-1} p \rangle$ denotes the **Hamiltonian** for the Riemannian metric g_q , we can show that **paths** $q(t)$ that **minimize length** locally and travel at **constant speed** follow a **coupled Ordinary Differential Equation** with a **momentum** vector $p(t)$:

$$\begin{cases} \dot{q}(t) = + \frac{\partial H}{\partial p}(q(t), p(t)) & \text{follow the velocity } v(t) = g_{q(t)}^{-1} p(t). \\ \dot{p}(t) = - \frac{\partial H}{\partial q}(q(t), p(t)) & \text{steer the momentum to stay on a geodesic path.} \end{cases}$$

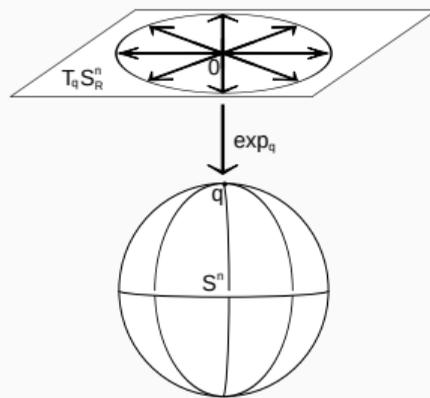
Geodesic paths are fully determined by:

- the starting **position** $q(t = 0)$,
- the starting **momentum** $p(t = 0) \iff$ the **velocity** $\dot{q}(t = 0) = g_{q(t=0)} p(t = 0)$.

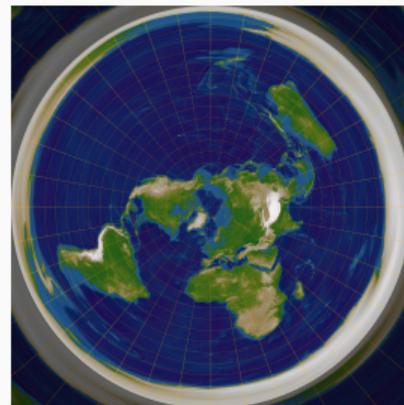
The exponential map [Chr10, Rok08]



From the 1D line
to the **circle**.



From the 2D plane
to the **sphere**.



Azimuthal equidistant
projection.

$\exp_{q_0}(p_0)$ denotes the **solution** $q(t=1)$ of the geodesic equation
with **initial condition** $q(t=0) = q_0, p(t=0) = p_0$.

Computing geodesics from A to B – three main scenarios

1. The metric has a lot of structure – **closed formulas**:

- Just like the Poincaré disk: simple example + symmetries.
- Riemannian **metric** \iff Relevant **kernel**.

2. The **metric** g_x is simple – path shortening:

- Discretize the path energy $\ell^2(\gamma) \simeq \frac{1}{N} \sum_{i=1}^N N^2 \|\gamma(i/N) - \gamma((i-1)/N)\|_{\gamma(i/N)}^2$.
- “Mean curvature flow”: gradient descent with respect to the snapshot positions.

3. The **cometric** $K_q = (g_q)^{-1}$ is simple – geodesic shooting:

- **Implement** the exponential $(q_0, p_0) \mapsto \exp_{q_0}(p_0)$ by integrating the Hamilton ODE.
- **Solve** the inverse problem $p_0 \mapsto \exp_A(p_0) \simeq B$ with an optimizer or a network.

Riemannian metrics provide:

- Expressive **vocabulary**: trees, balls, shapes and probability distributions.
- Complete **toolbox**: local metric \rightarrow geodesics, exponentials and barycenters.
- Appealing **message**:
simple paths on a curved space $>$ **complex paths on a flat space**.

This framework is the cornerstone of several applied fields,
and provides an **inspiring** outlook in many other settings.

\Rightarrow **Lab session** with GeomStats. \Leftarrow

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